# Conservation laws on a star-shaped network

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We consider a junction consisting of *m* incoming and *n* outgoing edges.



- Incoming edges:  $x \in \Omega_i = \mathbb{R}_-, i = 1, \dots, m$ ;
- Outgoing edges:  $x \in \Omega_j = \mathbb{R}_+, j = m + 1, \dots, m + n$ ;
- The junction is located at *x* = 0.

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On each edge we consider the evolution problem

$$\partial_t \rho_h + \partial_x f_h(\rho_h) = 0, \qquad h = 1, \dots, m + n,$$

• 
$$\rho_h$$
 conserved quantity,



•  $f_h : [0, R] \rightarrow \mathbb{R}_+$ , Lipschitz continuous,

• 
$$f_h(0) = 0 = f_h(R)$$
,

•  $\exists \bar{\rho} \in [0, R]$ , such that  $f'_h(\rho)(\bar{\rho} - \rho) > 0$ , for a.e.  $\rho \in [0, R]$ .

We postulate conservation at the junction

$$\frac{d}{dt}\sum_{h=1}^{m+n}\int_{\Omega_h}\rho_h(t,x)dx=0,$$

which we rewrite as

$$\sum_{i=1}^{m} f_i(\rho_i(t,0^-)) = \sum_{j=m+1}^{m+n} f_j(\rho_j(t,0^+)).$$



 Weak solutions, edge-wise entropy admissible

We call weak solution on the star-shaped network  $\vec{\rho} = (\rho_1, \dots, \rho_{m+n})$ 

• 
$$\rho_h \in L^{\infty}(\mathbb{R}_+ \times \Omega_h; [0, R]);$$

•  $\rho_h$  is a Kruzhkov entropy solution in  $\mathbb{R}_+ \times \{\Omega_h \setminus \partial \Omega_h\}$ . Namely  $\forall k \in [0, R]$  and  $\forall \varphi \in \mathcal{C}^1_c(\mathbb{R}_+ \times \Omega_h), \varphi \ge 0$ 

$$egin{aligned} &\int_{\mathbb{R}_+} \int_{\Omega_h} |
ho_h - k| arphi_t + ext{sign}(
ho_h - k) \left( f_h(
ho_h) - f_h(k) 
ight) arphi_x \, dx \, dt \ &+ \int_{\Omega_h} |u_0^h(x) - k| arphi(0,x) \, dx \geq 0 \ ; \end{aligned}$$

conservation at the junction holds.

#### >> Weak solutions are not unique in general.

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# The junction as a family of IBVPs

Fix 
$$\vec{u}_0 = (u_0^1, \dots, u_0^{m+n})$$
  
We look for  $\vec{\rho} = (\rho_1, \dots, \rho_{m+n})$  s.t.  $\forall h, \rho_h \in L^{\infty}(\mathbb{R}_+ \times \Omega_h, [0, R])$  solves

$$\begin{cases} \partial_t \rho_h + \partial_x f_h(\rho_h) = \mathbf{0}, & \text{on } ]\mathbf{0}, T[\times \Omega_h, \\ \rho_h(t, \mathbf{0}) = \frac{\mathbf{v}_h(t)}{\mathbf{v}_h(t)}, & \text{on } ]\mathbf{0}, T[, \\ \rho_h(\mathbf{0}, \mathbf{x}) = u_0^h(\mathbf{x}), & \text{on } \Omega_h, \end{cases}$$

where  $\vec{v} : \mathbb{R}_+ \to [0, R]^{m+n}$  is to be fixed at each t > 0

- to ensure conservation,
- depending on the state of the system,
- $\rightarrow$  encoding coupling conditions at x = 0.

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# ... Solves? Weak entropy solution for the IBVP

u is a weak entropy solution for the IBVP

$$\begin{cases} \partial_t u + \partial_x f(u) = 0, & \text{for } (t, x) \text{ in } \mathbb{R}_+ \times \mathbb{R}_- \\ u(t, 0^-) = u_b(t), \\ u(0, x) = u_0(x), \end{cases}$$

if

- *u* is a Kruzhkov entropy solution in the interior of  $\mathbb{R}_+ \times \mathbb{R}_-$ ,
- u satisfies the boundary condition in the sense of Bardos-LeRoux-Nédélec

$$sign(u(t,0^{-}) - u_b(t))(f(u(t,0^{-})) - f(k)) \ge 0,$$
  
 $\forall k \in \mathcal{I}(u(t,0^{-}), u_b(t)),$ 

which also write as

$$f(u(t,0^{-})) = \text{God}(u(t,0^{-}), u_b(t)).$$

Consider branch-wise constant data  $\vec{\rho} = (\rho_1, \dots, \rho_{m+n})$ 

To describe the Riemann Solver at the junction we define

► for 
$$i = 1, ..., m$$
  
Demand function :  $\Delta_i(\rho_i) = \max_s \operatorname{God}_{f_i}(\rho_i, s)$ ;

and we use them to determine the passing flow at the junction from each of the incoming roads

$$\Gamma_i: [0, R]^{m+n} \to [0, f_i^{\max}], \qquad i = 1, \ldots, m.$$

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### At a 1-2 divide

Fix a distribution factor  $\beta \in (0, 1)$ .

The passing flow at the junction is  $\Gamma_1 : [0, R]^3 \rightarrow [0, f_1^{max}]$  such that :

If βΔ<sub>1</sub>(ρ<sub>1</sub>) ≤ Σ<sub>2</sub>(ρ<sub>2</sub>), and (1 − β)Δ<sub>1</sub>(ρ<sub>1</sub>) ≤ Σ<sub>3</sub>(ρ<sub>3</sub>) then Γ<sub>1</sub>(ρ) = Δ<sub>1</sub>(ρ<sub>1</sub>),
 otherwise, Γ<sub>1</sub>(ρ) = min{β<sup>-1</sup>Σ<sub>2</sub>(ρ<sub>2</sub>), (1 − β)<sup>-1</sup>Σ<sub>3</sub>(ρ<sub>3</sub>)}.

In both cases

$$\begin{aligned} \mathbf{v}_{1} &= \left(f_{1|_{[\bar{\rho}_{1},\bar{P}_{1}]}}\right)^{-1}(\Gamma_{1}), \\ \mathbf{v}_{2} &= \left(f_{2|_{[0,\bar{\rho}_{2}]}}\right)^{-1}(\beta\Gamma_{1}), \\ \mathbf{v}_{3} &= \left(f_{3|_{[0,\bar{\rho}_{3}]}}\right)^{-1}((1-\beta)\Gamma_{1}). \end{aligned}$$

#### Remark

The application  $\vec{\rho} \mapsto (\Gamma_1, -\beta\Gamma_1, -(1-\beta)\Gamma_1)$  is not monotone.

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# At a 2-2 junction

We introduce a distribution matrix of the form

$$\mathbf{A} = \begin{pmatrix} \beta & \gamma \\ \mathbf{1} - \beta & \mathbf{1} - \gamma \end{pmatrix}$$

with  $\beta$  and  $\gamma$  in ]0,1[\{1/2}. Then

- $(\Gamma_1,\Gamma_2)\in [0,\Delta_1]\times [0,\Delta_2];$
- $A \cdot (\Gamma_1, \Gamma_2)^T$  must be in  $[0, \Sigma_3] \times [0, \Sigma_4]$ ;
- $\Gamma_1 + \Gamma_2$  should be as large as possible, under the constraints above.

### Remark

Counterexemples show that this solver lacks  $L^1$ -Lipschitz continuity with respect to the initial conditions.

See [Coclite-Garavello-Piccoli, 2005] and the book by Garavello and Piccoli Traffic Flow on Networks.

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# Vanishing viscosity approximations

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[Coclite-Garavello, 2010]
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# Fix $\varepsilon > 0$ and consider

$$\begin{cases} \partial_t \rho_h^{\varepsilon} + \partial_x f_h(\rho_h^{\varepsilon}) = \varepsilon \partial_{xx}^2 \rho_h^{\varepsilon}, \\ \sum_{i=1}^m (f_i(\rho_i^{\varepsilon}(t,0)) - \varepsilon \partial_x \rho_i^{\varepsilon}(t,0)) = \sum_{j=m+1}^{m+n} \left( f_j(\rho_j^{\varepsilon}(t,0)) - \varepsilon \partial_x \rho_j^{\varepsilon}(t,0) \right), \\ \rho_h^{\varepsilon}(t,0) = \rho_{h'}^{\varepsilon}(t,0), \\ \rho_h^{\varepsilon}(0,x) = u_{h,\varepsilon}^0(x), \end{cases}$$

where the initial conditions  $\vec{u}_{0,\varepsilon}$  approximates  $\vec{\rho}_0$ 

$$\begin{split} & u_{h,\varepsilon}^{0} \in W^{2,1} \cap C^{\infty}(\Omega_{h}; [0, R]), \\ & u_{h,\varepsilon}^{0} \longrightarrow \rho_{0,h}, \text{ a.e. and in } L^{p}(\Omega_{h}), 1 \leq p < \infty, \text{ as } \varepsilon \to 0, \\ & \left\| u_{h,\varepsilon}^{0} \right\|_{L^{1}(\Omega_{h})} \leq \| \rho_{0,h} \|_{L^{1}(\Omega_{h})}, \quad \left\| \partial_{x} u_{h,\varepsilon}^{0} \right\|_{L^{1}(\Omega_{h})} \leq TV(\rho_{0,h}), \quad \varepsilon \left\| \partial_{xx}^{2} u_{h,\varepsilon}^{0} \right\|_{L^{1}(\Omega_{h})} \leq C_{0}, \end{split}$$

with  $C_0 > 0$  independent from  $\varepsilon$ , *h*.

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For any fixed  $\varepsilon > 0$  there exists a unique  $\vec{\rho^{\varepsilon}}$  s.t.

$$\begin{split} \rho_h^{\varepsilon} &\in \mathcal{C}([0,\infty); \mathcal{L}^2(\Omega_h)) \cap \mathcal{L}_{loc}^1((0,\infty); \mathcal{W}^{2,1}(\Omega_h)), \quad \forall h, \\ 0 &\leq \rho_h^{\varepsilon} \leq \mathcal{R}, \qquad \sum_{h=1}^{m+n} \|\rho_h^{\varepsilon}(t,\cdot)\|_{\mathcal{L}^1(\Omega_h)} \leq \sum_{h=1}^{m+n} \|\rho_{0,h}\|_{\mathcal{L}^1(\Omega_h)}, \quad \forall t \geq 0, \end{split}$$

+ additional a priori estimates.

Compensated compactness  $\Rightarrow$  existence of a sequence  $\{\varepsilon_\ell\}_{\ell \in \mathbb{N}}, \varepsilon_\ell \to 0$ and a weak solution  $\vec{\rho}$  of the inviscid Cauchy problem at the junction s.t.

$$\rho_h^{\varepsilon_\ell} \longrightarrow \rho_h$$
, a.e. and in  $L_{loc}^p(\mathbb{R}_+ \times \Omega_h)$ ,  $1 \le p < \infty$ ,

for every  $h \in \{1, ..., m + n\}$ .

In [Andreianov-D.-Coclite, 2017] we further characterize the limit solution and prove its uniqueness. More details in the following...

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Analogously to [Diehl, 2009], [Andreianov-Mitrović, 2015] for m = n = 1

The condition  $\rho_h^{\varepsilon}(t,0) = \rho_{h'}^{\varepsilon}(t,0), \forall h, h' \in \{1,\ldots,m+n\}$ , translates into

 $\mathbf{v}_h(t)=\mathbf{v}_{h'}(t),$ 

for the family of hyperbolic IBVPs at the junction.

 $\vec{\rho} = (\rho_1, \dots, \rho_{m+n})$  is an admissible solution if there exists v in  $L^{\infty}(\mathbb{R}_+, [0, R])$  s.t.

- $\vec{\rho}$  is a weak solution,
- each component ρ<sub>h</sub> is weak entropy solution for the IBVP

$$\begin{cases} \rho_{h,t} + f_h(\rho_h)_x = 0, & \text{on } ]0, T[\times \Omega_h, \\ \rho_h(t,0) = v(t), & \text{on } ]0, T[, \\ \rho_h(0,x) = \rho_0^h(x), & \text{on } \Omega_h. \end{cases}$$

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We call germ of vanishing viscosity the set

 $\mathcal{G}_{VV} = \left\{ \vec{k} \in [0, R]^{m+n}, \text{ stationary edge-wise constant admissible solution} 
ight\}$ 

#### Lemma

If  $\rho_h$  is a Kruzhkov entropy solution in the interior of  $\mathbb{R}_+ \times \Omega_h$ ,  $\forall h \in \{1, \dots, m+n\}$ , TFAE

- $\vec{\rho}$  is an admissible solution;
- for a.e.  $t \in \mathbb{R}_+$ , the vector of traces  $\gamma \rho(t) = (\rho_1(t, 0^-), \dots, \rho_{m+n}(t, 0^+))$  is in  $\mathcal{G}_{VV}$ ;
- $\forall \vec{k} \in \mathcal{G}_{VV}, \vec{\rho} \text{ satisfies adapted entropy inequality on the network:}$  $<math>\forall \xi \in \mathcal{C}^{\infty}_{c}(\mathbb{R}_{+} \times \mathbb{R}), \xi \geq 0,$

$$\sum_{h=1}^{m+n} \left( \int_{\mathbb{R}_+} \int_{\Omega_h} \left\{ |\rho_h - k_h| \xi_t + \operatorname{sign}(\rho_h - k_h) (f_h(\rho_h) - f_h(k_h)) \xi_x \right\} \, dx \, dt \right) \geq 0.$$

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# Well-posedness for admissible solutions

### Theorem

- For any  $\vec{p_0}$  there exists an admissible solution  $\vec{p}$ .
- If  $\vec{\rho}$  and  $\vec{\rho}^*$  are admissible solutions corresponding to  $\vec{u}_0$  and  $\vec{v}_0$ , then

$$\sum_{h=1}^{m+n} \|\rho_h(t) - \rho_h^*(t)\|_{L^1(\Omega_h;\mathbb{R})} \le \sum_{h=1}^{m+n} \left\| u_h^0 - v_h^0 \right\|_{L^1(\Omega_h;\mathbb{R})}$$

# Fundamental properties of $\mathcal{G}_{VV}$

- completeness : we can associate an admissible solution to any Riemann datum.
- dissipativity : for any  $\vec{k}_1, \vec{k}_2$  in  $\mathcal{G}_{VV}$  with  $\vec{k}_\ell = (k_1^\ell, \dots, k_{m+n}^\ell), \ell = 1, 2,$

$$\sum_{i=1}^{m} \operatorname{sign}(k_{i}^{1}-k_{i}^{2}) \left(f_{i}(k_{i}^{1})-f_{i}(k_{i}^{2})\right) - \sum_{j=m+1}^{m+n} \operatorname{sign}(k_{j}^{1}-k_{j}^{2}) \left(f_{j}(k_{j}^{1})-f_{j}(k_{j}^{2})\right) \geq 0.$$

• maximality : if  $\vec{k}_1$  satisfies  $\uparrow$  for all  $\vec{k}_2$  in  $\mathcal{G}_{VV}$ , then  $\vec{k}_1 \in \mathcal{G}_{VV}$ .

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# Vanishing viscosity with different coupling conditions?

See [Guarguaglini-Natalini, 2015 & 2021] for the linear case

We consider coupling conditions inspired by the Kedem-Katchalsky conditions for membrane permeability

$$\begin{cases} \partial_t \rho_h^{\varepsilon} + \partial_x f_h(\rho_h^{\varepsilon}) = \varepsilon \partial_{xx}^2 \rho_h^{\varepsilon}, & t > 0, x \in \Omega_h, \\ \rho_h^{\varepsilon}(0, x) = \rho_{h,0}^{\varepsilon}(x), & h = 1, \dots, m+n, \\ f_i(\rho_i^{\varepsilon}(t, 0)) - \varepsilon \partial_x \rho_i^{\varepsilon}(t, 0) = \sum_j \mathfrak{c}_{ij}(\rho_i^{\varepsilon}(t, 0) - \rho_j^{\varepsilon}(t, 0)), & i = 1, \dots, m, \\ f_j(\rho_j^{\varepsilon}(t, 0)) - \varepsilon \partial_x \rho_j^{\varepsilon}(t, 0) = \sum_i \mathfrak{c}_{ij}(\rho_i^{\varepsilon}(t, 0) - \rho_j^{\varepsilon}(t, 0)), & j = m+1, \dots, m+n, \end{cases}$$

where  $c_{ij} > 0$ . We do not impose continuity at x = 0.

We can prove [Coclite-D. 2020]:

- Existence of parabolic approximations for any ε;
- Convergence (up to a subsequence) to a weak solution.

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# Characterization of the hyperbolic limit?

The 1-1 case

Assume  $f_1(\rho_{\chi}) = f_2(\rho_{\chi})$  and  $f_1(\hat{\rho}) = f_2(\check{\rho}) = \mathfrak{c}(\hat{\rho} - \check{\rho})$ 



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$$\begin{cases} \partial_t \rho_1^{\varepsilon} + \partial_x f_1(\rho_1^{\varepsilon}) = \varepsilon \partial_{xx}^2 \rho_1^{\varepsilon}, & t > 0, \ x < 0, \\ \partial_t \rho_2^{\varepsilon} + \partial_x f_2(\rho_2^{\varepsilon}) = \varepsilon \partial_{xx}^2 \rho_2^{\varepsilon}, & t > 0, \\ f_1(\rho_1^{\varepsilon}(t,0)) - \varepsilon \partial_x \rho_1^{\varepsilon}(t,0) = \mathfrak{c}(\rho_1^{\varepsilon}(t,0) - \rho_2^{\varepsilon}(t,0)), & t > 0, \\ f_2(\rho_2^{\varepsilon}(t,0)) - \varepsilon \partial_x \rho_2^{\varepsilon}(t,0) = \mathfrak{c}(\rho_1^{\varepsilon}(t,0) - \rho_2^{\varepsilon}(t,0)), & t > 0, \\ \rho_1^{\varepsilon}(0,x) = \hat{\rho}, & x < 0, \\ \rho_2^{\varepsilon}(0,x) = \check{\rho}, & x > 0. \end{cases}$$

As  $\varepsilon \to 0$  the limit is  $\rho_1(t, x) \equiv \hat{\rho}, \qquad \rho_2(t, x) \equiv \check{\rho}.$ 

- The couple  $(\hat{\rho}, \check{\rho})$  is a connection as introduced by [Adimurthi-Mishra-Gowda, 2005].
- Already obtained by adapted vanishing viscosity regularization

$$\begin{cases} \partial_t \rho_1^{\varepsilon} + \partial_x f_1(\rho_1^{\varepsilon}) = \varepsilon \partial_{xx}^2 a_1(\rho_1^{\varepsilon}), & t > 0, \ x < 0, \\ \partial_t \rho_2^{\varepsilon} + \partial_x f_2(\rho_2^{\varepsilon}) = \varepsilon \partial_{xx}^2 a_2(\rho_2^{\varepsilon}), & t > 0, \ x > 0, \end{cases}$$

where  $a_1$  and  $a_2$ :  $[0, 1] \rightarrow [0, 1]$  are strictly monotone increasing bijections and  $a_1(\hat{\rho}) = a_2(\check{\rho})$ .

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# Can be find a complete, maximal, $L^1$ dissipative germ for the limit problem?

 $\mathcal{G}_{KK} = \{(u_L, u_R) \text{ s.t. } u(t, x) = u_L \mathbb{1}_{\mathbb{R}_-} + u_R \mathbb{1}_{\mathbb{R}_+} \text{ is a stationary admissible solution}\}$ 

 $A = (\hat{\rho}, \check{\rho})$  defined by  $f_1(\hat{\rho}) = f_2(\check{\rho}) = \mathfrak{c}(\hat{\rho} - \check{\rho})$  must be in, together with  $u(t, x) \equiv 0$  and  $u(t, x) \equiv 1$ .

Of course, the germ contains all the couples (a, b) which

⇒ are traces of IBVPs with boundary conditions  $\vec{v} = (\hat{\rho}, \check{\rho}), \vec{v} = (0, 0)$  or  $\vec{v} = (1, 1)$ ;

➤ satisfy the Rankine-Hugoniot condition.

$$\mathcal{H} = \begin{cases} (a,b) \colon a \in [0,\check{u}_L] \cup \{\hat{\rho}\}, \\ b \in [\hat{u}_R, 1] \cup \{\check{\rho}\}, \\ f_1(a) = f_2(b). \end{cases} \cup \{(0,0)\} \cup \{(1,1)\}, \end{cases}$$

where  $\check{u}_L$  and  $\hat{u}_R$  satisfy  $f_1(\check{u}_L) = f_1(\hat{\rho})$  and  $f_2(\hat{u}_R) = f_2(\check{\rho})$ 

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Some Riemann problems do not have solutions if we impose that the traces at x = 0 of the solution are in  $\mathcal{H}$ .

We call  $\mathcal{G}_A^*$  the set of all couples (a, b) such that  $f_1(a) = f_2(b)$  and

 $\operatorname{sign}(a-\hat{\rho})\left(f_1(a)-f_1(\hat{\rho})\right)-\operatorname{sign}(b-\check{\rho})\left(f_2(b)-f_2(\check{\rho})\right)\geq 0,$ 

This set need to be in the germ because we need maximality and  $L^1$ -dissipativity.

A case by case study show that  $\mathcal{G}_{\mathcal{A}}^*$  is complete and maximal.

Also, we can show that each couple in  $\mathcal{G}_A^*$  correspond to the limit of a viscous profile.

Therefore, we can use it to prove well-posedness of limit solutions to the hyperbolic problem.



# Thank you for your attention!

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